Performance Evaluation of Airframe-Integrated Aerospike Propulsion Systems in Off-Design Flight Conditions

Hidemi Takahashi¹

Japan Aerospace Exploration Agency (JAXA), Kakuda, Miyagi, 981-1525, Japan

The thrust performance of airframe-integrated linear aerospike propulsion systems for high-speed aircraft operated under off-design in-flight conditions was evaluated, based on performance prediction models established for linear aerospike nozzles by accounting for the freestream effect. The airframe configuration, equipped with a boattail appearing to the right upstream of the primary cell nozzle exit, enabled the freestream effect to be utilized for increasing the spike surface pressure. Performances of each thrust component were predicted by using the physics-based models established in separately performed studies. First, a parametric study was conducted on a rocket primary thruster with various design parameters. Then, an advanced concept with an airbreathing primary propulsion system was evaluated for its feasibility by comparing its thrust performance with that of a conventional rocket primary thruster. The benefit of the aerospike nozzle integrated into the airframe was also evaluated by comparing its thrust to a configuration without the aerospike nozzle. Results showed that the employment of the aerospike nozzle offered improved performance from takeoff to cruise phases. The advanced concept using the bypass airbreathing system for the primary thruster followed by the aerospike configuration exhibited further improvement in thrust by compensating for the boattail drag.

Nomenclature

 $A = \text{area, m}^2$

CF = thrust coefficient

D = drag, N

M = Mach number

¹ Researcher, Engine System Research and Development Group, Space Transportation Mission Directorate, 1 Kimigaya-Aza-Koganezawa, Kakuda, Member AIAA.

 \dot{m} = mass flow rate, kg/s

P = pressure, kPa

 P_{0r} = total pressure of the cell nozzle flow, kPa

 P_r = static pressure at the cell nozzle exit, kPa

 P_w = static pressure on the nozzle surface, kPa

Re = Reynolds number

T = temperature, K

V = velocity, m/s

 $W = \text{wet area, m}^2$

x = coordinate in streamwise direction

y = coordinate in vertical direction

z = coordinate in spanwise direction

 γ = specific heat ratio

 ε = expansion ratio

 θ = angle, rad

Subscripts

0 = stagnation condition

a = freestream condition

b = ambient condition

boattail = boattail

bypass = bypass flow

c = combustion chamber

cell = cell nozzle

cb = cell base

cell_exit = cell nozzle exit

core = core flow

LE = leading edge

local = local contour

r = cell nozzle exit condition

r2 = bypass flow condition

sb = spike base

sf = skin friction force

spike = spike surface

straight = straight section

throat = throat

total = total

I. Introduction

Long-range hypersonic air transportation systems cruise over a wide range of altitudes and distances in the dense atmosphere after taking off horizontally from the ground; therefore, their flight envelopes involve off-design conditions for the propulsive nozzle, which is exposed to the supersonic external flow at various Mach numbers. The external flow surrounding the vehicle aftbody during flight considerably affects the performance of the nozzle owing to local changes in the ambient pressure and the resultant changes in flow properties of cell nozzle exhausts.

The linear aerospike nozzle is one of the most promising nozzle configurations for this kind of hypersonic vehicle because of its altitude compensating feature [1–3]. An engine employing an aerospike nozzle is capable of maintaining nearly optimal thrust performance from sea level to orbit, unlike that using a conventional bell-shaped nozzle, which reaches maximum performance only at one specific altitude (design altitude). Considering the needs of the vehicle during operation, such as thrust vectoring [4], and configurations that efficiently utilize the entire back area of a rectangular vehicle body to minimize the base drag, a clustered linear aerospike nozzle concept is beneficial.

When the aerospike nozzle concept is considered for deployment, it is necessary for mission feasibility to evaluate the thrust performance for given configurations. To date, numerous studies have been conducted to establish a performance prediction model, including a multidisciplinary optimization and feasibility study accounting for both aerodynamic and structural designs [1], and thrust performance and flowfield modelings based on experimental [5–7] and numerical simulation results [8-11]. Here, the overall thrust performance of a clustered linear aerospike nozzle composed with a rocket primary

thruster in the still-air condition (i.e., at sea level or in near-vacuum) can be expressed in terms of the aerodynamics with possible drags as follows [12]:

$$CF_{total} = \left(C_{cell} + \varepsilon_{cell} \times \frac{P_{cell_exit}}{P_c}\right) + \frac{P_{cb} \times A_{cb}}{P_c \times A_t} + \frac{\int P_{spike} \times Wdz}{P_c \times A_t} + \frac{P_{sb} \times A_{sb}}{P_c \times A_t} - \frac{D_{sf}[Pa]}{P_c \times A_t} \times W - \varepsilon_{spike} \times \frac{P_b}{P_c}$$

The left-hand side indicates the thrust coefficient, where the thrust value is normalized by the throat area and the combustion chamber pressure. On the right-hand side, the first two terms in parentheses represent the cell nozzle characteristics. The third term represents the cell base wake characteristics, which appear when a primary thruster comprises clustered cell nozzles. The fourth term is the thrust generated by the spike surface pressure, and the fifth term represents the spike base wake characteristics. Finally, the sixth and seventh terms are skin frictional drag and ambient pressure drag, respectively. The thrust equation can be divided into two portions: the primary inner thrust, comprising the first three terms, and the outer thrust comprising the remaining terms, including drags. The engine, which provides the primary inner thrust, can be either a rocket or an airbreathing propulsion system, depending on the mission and vehicle configurations. For the case of a vehicle during in-flight conditions, the thrust performance prediction model has to be modified by accounting for the influence of the external flow on each term.

To date, many insights on the influence of external flow have been provided through numerous studies examining the freestream effect [3,6,13–15], the mechanism of closed/open wake conditions [7], and clustering effects [8] in ground tests before flight experiments. It should be noted that actual flight experiments, including the Linear Aerospike SR-71 Experiment as a part of the X-33 program [16–17], have been conducted to investigate the interactions between the engine plume and external flow and to evaluate the methodologies to reduce base drags. Although the actual in-flight hot-flow test was not performed because of the concern of oxygen leakage in the test apparatus, airborne cryogenic gas flow pluming tests have been conducted, and the expected data as well as the data for the aerodynamic flight performance were obtained [18]. Nasuti et al. [3] suggested that the shroud base pressure governs the main mechanism of flow expansion and base wake transition characteristics. In addition to this insight, previous studies [12,14-15] investigated the mechanism behind the spike surface pressure generated by the external flow coupled with clustered cell jet flows, base characteristics of the cell and spike, and skin friction characteristics in the presence of the external flow. The results showed that the thrust performance can be significantly improved under off-design in-flight conditions in the presence of external flow with a certain aftbody configuration, which appears as an integrated segment in the airframe. This insight is particularly important because the flight

performance was expected to decrease under transonic flight conditions [13,16] and off-design flight conditions [19] because of locally decreased flow static pressure acting on the nozzle surface and flow instabilities [16]. Additionally, a large angle of the boattail may lead to pressure reduction on its surface, and thus, generate drag. To account for these insights and concerns, it is important to evaluate the performance of the aerospike nozzle in the presence of external flow and to evaluate the effectiveness of the boattail configuration. To this end, it is vital to update the thrust performance prediction model for application to in-flight conditions. Additionally, it is beneficial to demonstrate the supersonic cruise efficiency by evaluating an airframe-integrated propulsion system equipped with an aerospike nozzle for a hypersonic air transportation system.

This study had two major goals: to evaluate the aerodynamic performance of a clustered linear aerospike nozzle operated under off-design in-flight conditions, and to conduct a conceptual study evaluating the effectiveness of an advanced airframe-integrated propulsion. While the thrust generated by the linear aerospike nozzle portion was considered as the outer thrust, the primary inner thrust was evaluated for the cases of a rocket or an airbreathing propulsion system (e.g., a ramjet). The performance evaluation for the baseline concept was based on the quasi-one-dimensional analytical models established in previous studies [12,14,15], which predict thrust value for each thrust component. Here, an analytical model denotes a physics-based model established without involving any empirical approximation. Although the quasi-one-dimensional analysis would limit the applicability of its results because the flowfield would be highly two- or three-dimensional, the analysis is useful for addressing the trend of performance for a new configuration. Additionally, it is worth noting that the analytical model established in the previous studies [14,15] indicated that it can predict the spike surface pressure for three-dimensional cases as well as two-dimensional cases with an uncertainty less than 5%. Since this study emphasizes the conceptual analysis of the airframe-integrated aerospike propulsion system and evaluating its feasibility, the vehicle aftbody was mainly considered. Additionally, based on the performance prediction models, advanced propulsion concepts that involve the bypass ramjet concept were evaluated for their thrust performance. Since this study focused on the performance evaluation, structural and weight impacts on the system feasibility was not discussed in this study.

II. Baseline Airframe and Nozzle Configurations and Flight Conditions

A. Airframe Configuration

Figure 1 illustrates an example configuration of the center cut of the vehicle body under supersonic flight condition with installation of the aerospike nozzle. Although Figure 1 illustrates the baseline case in which a rocket propulsion system is

considered for its primary thruster [12,14,15], this study emphasizes a cruise vehicle with an aerospike nozzle. The airframe has a leading edge, the angle of which will be varied, and the dimensions were determined by other design parameters as discussed below. A boattail appears to the right upstream of the primary thruster exit. Because the boattail portion is inclined at a certain angle relative to the center axis of the airframe, the Mach number of the boattail can be different from that of the flight; it can be greater than that of the external flow because of flow expansion. Additional design parameters, which will be described in a later section, were consistent with those in previous studies. For the advanced propulsion concept, the primary thruster portion would be replaced with an airbreathing propulsion system. These details will be discussed in a later section, along with an evaluation of the thrust performance. The overall airframe geometry and dimensions were referred from a precursor study [20], and some modifications were considered in the aftbody portion to install the aerospike nozzle with a rocket primary thruster. The overall airframe length was considered to be 40 m [20].

B. Linear Aerospike Nozzle Configuration (Airframe External Nozzle)

The aerospike nozzle portion comprises a nonclustered two-dimensional cell nozzle (designated as 2-D nozzle) or clustered cell nozzles with a designed Mach number of 3.5 [12,14,15], followed by a straight planar section and a contoured spike section. In the case that the cell nozzle is clustered, the cell base appears between two neighboring cell nozzle jets. The straight section appears because the Mach number of the cell nozzle is 3.5 rather than unity, and therefore, the Mach wave emanating from the cell nozzle lip impinges on the nozzle wall; behind the Mach wave impingement point (the end point of the straight section), the contour was designed using the method of characteristics with given parameters, which will be described below, to eliminate reflected expansion waves. In the case that the spike nozzle is truncated, the spike base appears. This feature is shown on the right side of Figure 1 as an end-view of the back surface of the rectangular vehicle. The width of the projected airframe back area was fixed at 6.0 m, and the height was determined according to the boattail angle and resultant spike geometry: the total height was first assumed to be 3.0 m, but the actual height was determined according to the given boattail angle and length, and resultant spike contour. The specific heat ratio (γ) of the cell nozzle exhaust gas was assumed to be 1.2, which is the approximate value of the specific heat ratio in the real engine exhaust gas flow, instead of the value of 1.4 for air. Note that the analytical model has been validated for the specific heat ratio of 1.4 by comparing the spike surface pressure to that obtained by the corresponding cold-flow experiment [12,14,15]. The analytical models used in this study were assumed to be applicable in the case of a specific heat ratio of 1.2 as well.

Geometries of the cell and spike nozzles were designed by using the method of characteristics depending on the given airframe geometries, cell nozzle exit Mach number (which determines the length of the straight section), specific heat ratio of the cell nozzle exhaust gas, and angle and length of the boattail. The baseline contour geometry of the spike is presented in Figure 2, which was calculated for the case of the inclination angle of the cell nozzle axis of 25°, boattail length of 2.0 m, and cell nozzle exit height of 0.48 m. The angle and dimension were chosen to match the experiment [15]. The boattail angle will be discussed in detail below. The length of the straight section is 1.61 m, the total length of the spike is 7.31 m, and the height is 1.69 m. The ideally expanded cell nozzle flow would have a Mach number of 4.7 at the end of the spike nozzle.

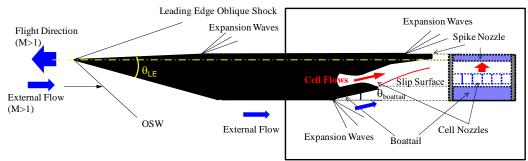


Figure 1 Illustration of baseline airframe configuration with boattail and aerospike nozzle.

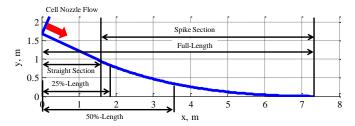


Figure 2 Baseline geometry of spike nozzle.

C. Flight Conditions

Figure 3 shows the flight conditions simulated in this study. The flight conditions were determined by using the nozzle pressure ratio (NPR), P_{0r}/P_b , which was defined as the pressure ratio between the cell nozzle total pressure (P_{0r}) and the environmental pressure (P_b). The NPR identified the jet expansion condition and simulated the flight altitude condition: a higher flight altitude condition corresponded to a larger NPR condition by assuming a constant combustion chamber pressure (P_{0r}) with varying environmental pressures (P_b). The designed NPR of the cell nozzle, which corresponded to its optimal expansion condition, was 121.3. Note that the specific heat ratio (γ) used for the cell nozzle exhaust gas was 1.2 as aforementioned. The flow structure of the cell nozzle exhaust flow formed downstream of the cell nozzle exit changed according to a given NPR condition. The relationship between the flow structure and NPR range is summarized in Table 1.

The nozzle flow separation condition was obtained by using the criterion suggested by Stark [21]. The expansion condition is summarized in terms of the cell nozzle NPR. The optimal expansion condition for the spike nozzle, which had the exit Mach number of 4.7 at the end of the spike nozzle, with the parameters described here is 1131.0. The simulated flight altitude for the cruise conditions is up to 24 km above sea level. This altitude condition may likely involve off-design flight conditions. The atmospheric pressure was calculated based on the U.S. Standard Atmosphere [22].

Figure 4 shows the relationship between the cruise Mach number, boattail Mach number, and leading edge angle. The vertical axis is the ratio between the calculated boattail Mach number and cruise Mach number for each value of the leading edge angle. The Mach number ratio ($M_{boattail}/M_{cruise}$) appears to be insensitive to the leading edge angle; for instance, its variation for the calculated leading edge angles is less than 0.7% at the cruise Mach number of 2.0.

Since the maximum boattail angle that limits the separation of the external flow on the boattail is not accurately known because of the viscous effects, the separation limit was determined by using the separation criterion [21] and the following relationship:

$$\frac{P_{boattail}}{P_{cruise}} > \frac{P_{separation}}{P_{cruise}} = \frac{\pi}{3 \cdot M_{separation}} \tag{1}$$

where $P_{separation}$ and $M_{separation}$ are the wall pressure and Mach number, respectively, where the flow separation starts. The calculated maximum boattail angle was approximately 15° up to the flight Mach number 4.7. This angle was also considered as the baseline boattail angle.

The assumed flight trajectory is presented in Figure 5. The flight profile considered here was a one-way flight path, which was divided into three phases: take-off, ascent, and cruise. The assumed flight Mach number required for take-off was 0.4. During the acceleration in the ascent to cruise phases, the flight Mach number was calculated by a fixed dynamic pressure (*q*) condition at 50 kPa [23]. The flight Mach number was considered up to 4.0, and therefore, the Mach number ceiling occurs at 4.0. For the baseline analysis, only the cruise condition was evaluated in detail to match the experiment [15] with the altitude of 13 km above sea level and flight Mach number of 2.0. The associated Reynolds number based on the airframe length is 2.6×10⁹ for the flight condition of Mach number 2.0 and the altitude of 13 km, where the ambient temperature and pressure are 216.6 K and 15 kPa, respectively [22]. Calculating the Reynolds number for the other flight conditions indicates that the flow around the airframe and on the spike surface is clearly in the turbulent regime, and therefore, the boundary layer condition around the airframe was considered to be turbulent.

The combustion chamber pressure, which corresponds to the total pressure of the cell nozzle, was fixed at 4.0 MPa [24]. A realistic combustion chamber pressure value was used here to enable a comparison between the calculated thrust performance and that of existing engines.

Table 1 Cell jet expansion features according to NPR range.

NPR range	Cell jet expansion condition	Note
NPR > 121.3	Underexpansion	-
NPR = 121.3	Optimal expansion	At designed NPR
38.6 < NPR < 121.3	Overexpansion	No nozzle separation involved
22.4 < NPR < 38.6	Overexpansion	Nozzle flow separation takes place
NPR < 22.4	Overexpansion	Normal shock wave exists inside nozzle

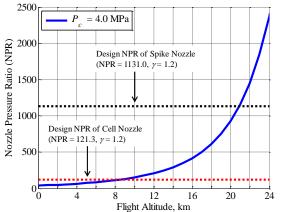


Figure 3 Simulated flight conditions: NPR against Flight altitude

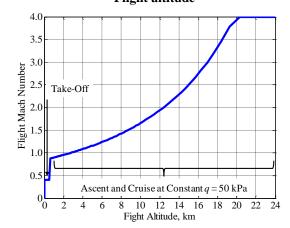


Figure 5 Hypothetical flight trajectory.

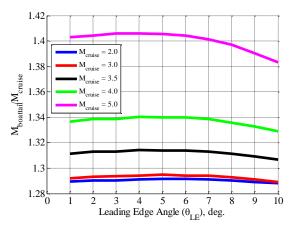


Figure 4 Relationship among leading edge angle, cruise Mach number and boattail Mach number.

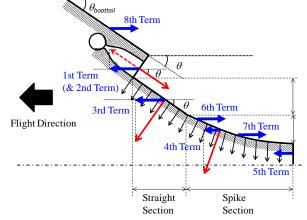


Figure 6 Pressure distributions acting on the nozzle surface.

III. Thrust Performance Prediction Model

A. Overall Thrust Performance Prediction Model Under In-Flight Conditions

A rocket thruster is considered as the primary thruster for the baseline configuration. The overall performance prediction model, based on quasi-one-dimensional analysis [15] in the case of a rocket primary thruster and boattail-equipped configuration (the configuration illustrated in Figure 1) under in-flight conditions, is expressed as Eq. (2) by accounting for the pressure distributions and their directions acting on the airframe and nozzle surfaces, as shown in Figure 6.

$$CF = \sqrt{\gamma_{r} \cdot \left(\frac{2}{\gamma_{r}+1}\right)^{\frac{\gamma_{r}-1}{\gamma_{r}-1}}} \cdot \left(\sqrt{\frac{2\gamma_{r}}{\gamma_{r}-1}} \left\{1 - \left(\frac{P_{r}}{P_{0r}}\right)^{\frac{\gamma_{r}-1}{\gamma_{r}}}\right\} + \left(\frac{P_{r}}{P_{0r}} - \frac{P_{b}}{P_{0r}}\right) / \left[\left(\frac{P_{r}}{P_{0r}}\right)^{\frac{1}{\gamma_{r}}} \sqrt{\frac{2\gamma_{r}}{\gamma_{r}-1}} \left\{1 - \left(\frac{P_{r}}{P_{0r}}\right)^{\frac{\gamma_{r}-1}{\gamma_{r}}}\right\}\right]\right) \cdot \cos\theta_{boattail} + \frac{\int P_{cb} \cdot dA_{cb}}{P_{0r} \cdot A_{throat}} \cdot \cos\theta_{boattail} + \frac{\int P_{straight} \cdot \sin\theta_{boattail} \cdot dA_{straight}}{P_{0r} \cdot A_{throat}} + \frac{\int P_{spike} \cdot \sin(\theta_{boattail} - \theta_{local}) \cdot dA_{spike}}{P_{0r} \cdot A_{throat}} + \frac{\int P_{sb} \cdot dA_{sb}}{P_{0r} \cdot A_{throat}} - \frac{\int D_{sf} \cdot \left\{\sin\theta_{boattail} \cdot dA_{straight} + \sin(\theta_{boattail} - \theta_{local}) \cdot dA_{spike}\right\}}{P_{0r} \cdot A_{throat}} - \frac{P_{b} \cdot \int \left\{\sin\theta_{boattail} \cdot dA_{straight} + \sin(\theta_{boattail} - \theta_{local}) \cdot dA_{spike} + dA_{sb}\right\}}{P_{0r} \cdot A_{throat}} - \frac{P_{boattail} \cdot A_{boattail} \cdot \cos\theta_{boattail}}{P_{0r} \cdot A_{throat}}$$

Here, the effects of change in environmental pressure (P_b) and total pressure (P_{0r}) are considered. The left-hand side is the coefficient of aerodynamic thrust performance. On the right-hand side, the first term corresponds to the thrust generated by the cell nozzle. The second term is the drag generated by the cell base that appears when the clustered cell nozzle configuration is deployed. The third and fourth terms are the thrusts generated from the straight and spike sections, respectively. The fifth term is the thrust generated from the spike base pressure. This could also be a drag, depending on the truncation configuration and flow condition. For the full-length spike model, the fifth term on the right-hand side is ignored. The sixth term is the skin friction drag, the seventh term is the ambient pressure drag, and the eighth term is the boattail drag. The local contour angle, θ_{local} , was taken into account when integrating the spike surface pressure. The boattail angle, $\theta_{locatail}$, was a variable parameter and other design parameters are summarized in Table 2. Prediction models for each thrust component are described in more detail in the following sections. The leading edge angle and length should be determined by accounting for a given design cruise Mach number and certain other design parameters for the inlet, but these are not discussed in detail in this study.

Table 2 Summary of Design Parameters.

Tuble 2 Summary of Design 1 arameters.		
Fixed Parameters	Value	
Combustion chamber pressure (P_{0r})	4.0 MPa	
Cell nozzle exit Mach number (M_r)	3.5	
Aftbody height	Determined by other parameters	
Aftbody width	6.0 m	
Spike length (truncation rate)	0–70%	
Flight altitude	0–24 km	
Number of cell nozzles	0 (2-D), 6 (One Side)	
Aspect ratio of cell nozzle exit	1.0	
Wall thickness	0.01 m	
Boattail angle	10–15°	
Boattail length	2.0 m	

B. Characteristics of Cell Nozzles

The cell nozzle characteristics were calculated from the first term in Eq. (2) based on the aerodynamic one-dimensional analysis. The exit dimension of each cell nozzle that has a rectangular shape was determined by using the number of cell nozzles, airframe base width (6.0 m), cell base width, aspect ratio of each cell nozzle (1.0), and wall thickness (0.01 m). The cell base width is preset to 0 for the case of a 2-D cell nozzle. Note that the clustered cell nozzle configuration is more realistic considering the benefits in operation as mentioned in the Introduction; however, the 2-D cell nozzle configuration was also simulated for a comparison. Because the exit Mach number of each cell nozzle was presumed to be 3.5 to enable a comparison between the simulated results and the experimental ones [14,15], the throat area was determined by the exit area of the cell nozzle and the expansion ratio, which was determined by the area-Mach number relationship.

C. Spike Surface Pressure Distributions

The spike surface pressure distributions in the presence of the external flow were predicted by the analytical model and procedure established in the previous study [15], as shown in Figure 7. As stated in Takahashi et al. [15], the external flow governs the pressure determination mechanism on the spike surface according to the physics for the formation of a slip line between the cell nozzle flow and the supersonic external flow, and the appearance of double oblique shock waves at the spike entrance. Additionally, the model is applicable to not only the 2-D cell nozzle configuration but also to the clustered cell nozzle configurations, regardless of the cell base width, with maximum error of approximately 5% [15]. More details about the procedure for deriving the spike surface pressure distribution are provided in Takahashi et al. [15]. It should be noted that a slight modification was introduced into the analytical model so that the model can be applicable for the case when the

boattail angle and cell nozzle inclination angle are different, because the spike surface plays a role in the compression wall for the boattail flow. When the difference in pressure between P_5 and P_6 in Figure 7 exceeded the convergence limit (i.e., 10% [15]), the averaged pressure between P_5 and P_6 was taken. This difference was mainly caused by the slight discrepancy between the assumed flow turning angle for the external flow and the local contour angle. The difference between these two pressures was considered the error value.

Figure 8 compares the calculated spike surface pressure distributions in the presence of external flow for various NPR conditions. Normalized pressure, which is the surface pressure divided by the cell nozzle total pressure, is plotted against the streamwise direction. The Mach number of the external flow was 2.0, as observed in the experiment. The boattail Mach number was 2.6, as seen in Figure 4, and the resulting static pressure of the external flow was 5.2 kPa. The error was determined from the difference between P_5 and P_6 , as mentioned above. As the cell nozzle total pressure increases, the overall spike surface pressure decreases. This is because the wall pressure changes depending on the cell nozzle expansion feature. For example, under overexpansion condition (low cell nozzle total pressure), the overall wall pressure is high because of the appearance of an oblique shock wave on the wall. In contrast, during the underexpansion condition (high cell nozzle total pressure), the overall spike surface pressure decreases because of the Prandtl-Meyer expansion formation on the spike surface. This same feature was observed in the experiment [15]; therefore, this model is expected to accurately predict the pressure distribution.

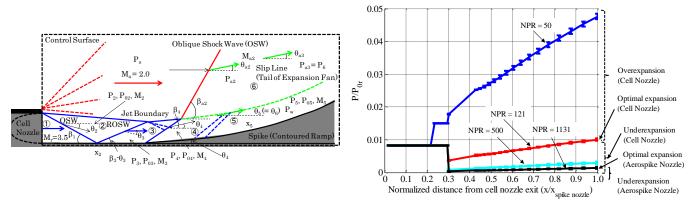


Figure 7 Schematic of flowfield model with external flow for an overexpanded jet.

Figure 8 Simulated spike surface pressure distributions for various cell nozzle total pressure cases.

D. Cell Base and Spike Base Pressures

The characteristics of the cell base and spike base pressures were calculated by using the base pressure prediction models established in previous studies [14–15]. Only a brief explanation and simulated results are presented in this section. Detailed descriptions of the model and the procedure for deriving the base pressure are available in Takahashi et al. [14–15].

The spike base pressure was predicted by using the model depicted in Figure 9. An overexpansion case is presented here. For the underexpanded jet, the oblique shock waves in regions 1–3 and the compression waves in region 5 are replaced with Mach waves. The premises for the model are that the spike base wake structure is in the closed-wake state and that the spike base pressure is highly correlated with the pressure at the spike edge, as observed in the experiment [15]. Additionally, the lip shock and recompression shock waves emerge behind the Prandtl-Meyer expansion fan formed to the center at the spike edge.

Figure 10 presents simulated spike base pressure distributions for the case in which the total pressure of the cell nozzle is 4.0 MPa. The normalized pressure in terms of the normalized differential pressure (NDP), defined as $(P_{sb} - P_b)/P_{0r}$, is plotted against the NPR. This normalization is useful when the pressure is compared to the ambient pressure level. In this calculation, the ambient pressure was varied, and thus, the spike base pressure was simulated according to the change in NPR. The spike length was truncated at 25% for this calculation. Under a condition of very low NPR, which involves an overexpanded cell nozzle, the calculated base pressure can be a positive value. In the NPR range up to approximately 120, the base pressure periodically drops and increases as the NPR increases. The hypothesis for this behavior is that the Prandtl-Meyer expansion fan enters the spike surface as the cell nozzle is underexpanded. Above NPR of approximately 120, the base pressure drops again, but then increases as NPR increases; eventually, the base pressure becomes higher than the ambient pressure. This is because the Prandtl-Meyer expansion fan emanating from the cell nozzle lip enters the spike base wake region, and consequently the base wake pressure is affected by the cell nozzle pressure condition. Thus, the base pressure basically acts as a drag or neutral thrust, except under conditions of NPR less than approximately 120 and greater than 1000.

Regarding the cell base pressure distribution, the preceding model was applied with the same assumptions because the model was established based on flow visualization results to more accurately represent flow physics.

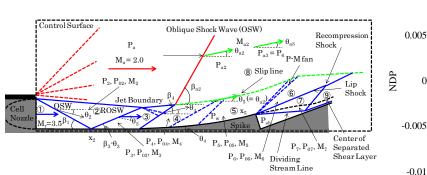


Figure 9 Schematic of the flowfield model with external flow for an overexpanded jet.

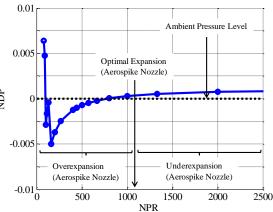


Figure 10 Simulated spike base pressure distribution plotted against NPR.

E. Skin Friction Drag

A previous study [12] showed that the wall impingement by oblique shock waves significantly increased the skin friction drag, whereas expansion fans negligibly affected the skin friction values. The skin friction drag was predicted by using the Van Driest II model [25]. This model can be applied when no separation or shock waves are involved. In the case that oblique shock waves exist, the skin friction drag drastically increases on the entire spike surface [12]. Because no analytical model has been established to accurately predict the increased skin friction values, the skin friction drag was assumed based on the Van Driest II model, which is given as follows:

$$\frac{\sin^{-1}\sqrt{\frac{\gamma-1}{2}M_{ae}^2\frac{T_e}{T_w}}}{\sqrt{C_f\left(\frac{T_w}{T_e}-1\right)}} \cong 4.15\log\left(\operatorname{Re}_{xe}C_f\frac{\mu_e}{\mu_w}\right) + 1.7$$
(3)

where M_{ae} is the Mach number of the flow at the location of interest on the spike surface, T_e is the static temperature of the flow. M_{ae} and T_e were calculated by the model described in Figure 7. The viscosities μ_e and μ_w for the free stream and at the wall, respectively, were calculated by the Sutherland viscosity law. The wall temperature (T_w) was assumed to be 1000 K because the temperature is expected to be close to that of the turbine disk [24]. As mentioned above, since the effect of shock wave impingement was not considered in this calculation, the skin friction was estimated by taking the mean value over the entire spike surface. The calculated mean skin friction coefficient was found to be 2.85×10^{-2} and skin friction drag was 92.1 kN for the baseline condition mentioned above.

F. Boattail Drag (Aftbody Drag)

The boattail drag was calculated with the eighth term in Eq. (2) by considering the angle and length of the boattail. Additionally, the external flow on the boattail is presumed to flow on the surface with no separation. The angle and length of the boattail were first determined as 15° and 2.0 m, respectively, in the baseline case; however, they were later treated as variables. The maximum boattail angle that may begin flow separation was assumed to be 15° as mentioned above. A more detailed investigation on the separation angle should be carried out for a more accurate prediction.

In addition to the boattail drag, this study did not account for the forebody drag that appears as the leading edge drag, although the leading edge drag has to be considered for accurately predicting the net thrust of the propulsion system. Note that methodologies to reduce the boattail drag can be found elsewhere [26].

IV. Results and Discussion for Baseline Configuration

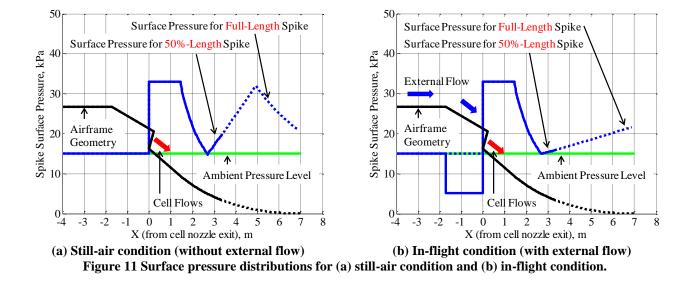
The insights gained from the performance evaluation for the airframe-integrated aerospike propulsion system are discussed in this section and cover the following topics of interest: pressure distributions on the airframe and nozzle wall surfaces, thrust components, and thrust performance with various parameters. Since this section emphasizes the propulsion system, the flight stability will be discussed in the next section.

A. Surface Pressure Distributions

Figure 11 presents pressure distributions of the airframe surface for (a) still-air condition and (b) in-flight condition at Mach 2.0. For the baseline case, the 2-D cell nozzle was considered with a full-length or 50%-length spike, $P_{0r} = 4.0$ MPa, and boattail angle of 15°. The baseline case was similar to that in the experiment [15]. Additionally, as aforementioned, only the nonreacting flowfield was considered for evaluating the aerodynamic performance. The ambient pressure level for still-air and in-flight conditions was fixed at 15 kPa, which corresponds to the conditions of flight altitude of approximately 13 km above sea level or flight Mach number of approximately 2.0 at sea level condition. The still-air condition was calculated for comparison.

For the still-air condition (Figure 11(a)), the surface pressure distribution was calculated by using the method of characteristics with consideration of the shock wave reflections on the spike surface. Pressures for both the full-length and 50%-length spike cases are plotted. Additionally, the airframe geometry is superimposed. The NPR for this condition was approximately 267; therefore, the exhaust plume from the cell nozzle (as an inner nozzle) is underexpanded; the aerospike nozzle (as an outer nozzle) is overexpanded. The shock wave reflection on the spike surface, which was observed for cases of both full- and 50%-length spikes, features the pressure distribution with periodic compression and expansion.

For the in-flight condition (Figure 11(b)), the spike surface pressure was calculated by using the analytical prediction model [15]. As previously stated in Takahashi et al. [15] in the case of in-flight conditions, unlike still-air conditions, the supersonic external flow causes the spike surface pressure to continue increasing as flow proceeds downstream.

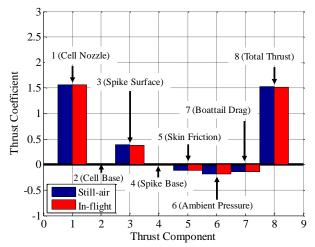


B. Thrust Components

Figure 12 shows thrust components in the case of a Mach 2.0 flight condition with the 50%-length spike configuration and boattail angle of 15°. The ambient pressure was fixed at 15 kPa, which corresponds to a supersonic cruise altitude of approximately 13 km above sea level. The case of a still-air condition is also plotted for comparison. The cell nozzle total pressure was 4.0 MPa. Each thrust component is numbered; the number corresponds to each term in Eq. (2). The cell nozzle thrust clearly makes the largest contribution to the total thrust for both still-air and in-flight conditions. The thrust generated by the spike surface pressure makes the second largest contribution to total thrust. Whereas the drags generated by the ambient pressure and the boattail portion of the airframe are comparable, the ambient pressure drag has the largest negative impact on the total thrust. The skin friction drag cannot be neglected as well. With this flight condition and configuration (i.e., spike length), there is no remarkable difference in thrust between the still-air and in-flight conditions.

Figure 13 shows the thrust components for the full-length spike case. Each value is plotted as in Figure 12. A remarkable difference in thrust between the still-air and in-flight conditions is clear for the spike surface pressure. Although the spike surface pressure under the in-flight condition was expected to dramatically decrease because of the existence of the external flow [13,16], the spike surface pressure remained comparable to that for the still air. As mentioned before, it is desirable to maintain spike surface pressure higher than the ambient pressure to generate thrust. The mechanism that maintained the spike surface pressure high was the formation of the oblique shock wave and slip line between the cell nozzle flows and external flow as described in previous sections. By comparing the thrust with that presented in Figure 12 and considering the trade-off between the spike length and the associated skin friction drag, the following statement can be made: the longer the spike, the

higher the thrust performance under in-flight conditions. Therefore, a longer spike length is preferable to maximize the spike surface pressure distributions, and thus, increase the thrust performance for in-flight conditions. It should, however, be noted that the skin friction drag may increase in the flowfield where strong impingement of shock waves is involved.



3 (Spike Surface) 8 (Total Thrust)

2.5 (Ambient Pressure)

1.5 (Skin Friction)

2 (Cell Base)

4 (Spike Base)

7 (Boattail Drag)

Thrust Component

Figure 12 Thrust components for still-air and in-flight conditions in the case of 50%-length spike.

Figure 13 Thrust components for still-air and in-flight conditions in the case of full-length spike.

C. Effects of Flight Conditions

Figure 14 shows the thrust performance calculated for various flight Mach numbers for cases with 50% and full-length spikes, 2-D cell nozzle, and boattail angle of 15°. The ambient pressure was fixed at 15 kPa, which corresponds to a supersonic cruise altitude of approximately 13 km above sea level. The thrust performance calculated for only the cell nozzle is additionally plotted for comparison.

The total thrust performance for the 50%-length spike could be lower than that of the cell nozzle for the Mach number above approximately 2.2. Additionally, the total thrust performance for the 50%-length spike shows that the performance decreases as flight Mach number increases, but the value asymptotically converges at approximately 1.5 at higher Mach numbers. In contrast, the total thrust performance for the full-length spike shows a similar trend to that of the 50%-length spike case, but much higher values over the calculated Mach number range. The thrust performance asymptotically converges at approximately 1.85 in the higher Mach number range. Thus, a longer spike length configuration significantly benefits total thrust, regardless of the flight Mach number.

D. Effects of Boattail Angles

Figure 15 shows the thrust performance calculated for various boattail angles for cases with 50% and full-length spikes, 2-D cell nozzle, and various boattail angles ranging from 10° to 20°. Note that the maximum boattail angle was stated in the previous section to be 15°, but the angles for additional 5° were also calculated for comparison. The cell nozzle total pressure, flight Mach number, and ambient pressure were fixed at 4.0 MPa, 2.0, and 15 kPa, respectively. The thrust generated by the cell nozzle itself is additionally plotted for comparison. The cell nozzle thrust remained constant as the inclination angle of the cell nozzle was kept constant at 25°. In contrast, for the both cases with the 50%-length and full-length spikes, the total thrust slightly decreased as the boattail angle increased because of the increasing boattail drag. Similar to the observation in Figure 14, the longer spike configuration tends to have higher thrust performance. The results indicated that the boattail angle is less sensitive to the total thrust for this configuration and conditions.

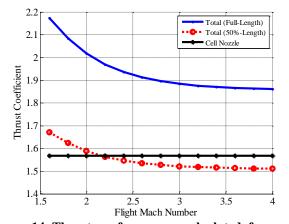


Figure 14 Thrust performances calculated for various flight Mach numbers at 50%- and full-length spikes.

Figure 15 Thrust performances calculated for various boattail angles at 50%- and full-length spikes.

E. Effects of Spike Truncation Length

Figure 16 presents the effect of the spike length on the total thrust performance. This figure plots the thrust performance calculated for flight Mach numbers of 1.5, 1.75, 2.0, 2.25, 2.5, 3.0, and 4.0, and for the boattail angle of 15° for various spike lengths. Similar to the calculations given for Figures 14 and 15, the cell nozzle total pressure and ambient pressure were fixed at 4.0 MPa and 15 kPa, respectively.

In general, the total thrust linearly increases with the spike length increasing toward the full length for all flight Mach number cases. A notable point is that the total thrust performance could be lower than that of the cell nozzle thrust for all flight Mach numbers with a spike length shorter than the length at a certain truncation ratio of approximately 35% to 58%

depending on the flight Mach number. The total thrust performances for the flight Mach numbers above approximately 2.0 become almost identical, and the thrust value with the spike length longer than 60% becomes higher than that of the cell nozzle.

In summary, it appears that a spike longer than 60% is beneficial for all flight Mach number range considered in this study. From the observation obtained in this section, it can be concluded that an optimization study of airframe geometry is possible for a given mission.

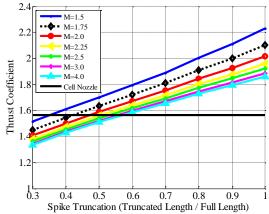


Figure 16 Thrust performances calculated for various spike truncation lengths.

V. Advanced Airframe-Integrated Propulsion Concepts

Based on the performance prediction models discussed above, advanced concepts of airframe-integrated propulsion systems are evaluated with respect to the thrust performance. The aerospike nozzle configuration in which a boattail is equipped to the right upstream of the primary cell thruster is considered. The mechanism behind the increase in thrust is, as aforementioned, attributed to the increase in the surface pressure of the spike under supersonic flight conditions [14–15]. However, a boattail-equipped airframe configuration with a sufficiently long boattail, which allows the external flow to flow on the boattail surface without separation and to be inclined toward the center axis of the airframe, may not always be allowed due to geometrical constraints of the airframe. Therefore, certain thrust augmentation (e.g., secondary flow) would be desirable to increase the spike surface pressure, and thus, increase the thrust. A feasibility study is conducted for advanced airframe-integrated propulsion systems, along with an evaluation of the thrust performance in a hypothetical flight trajectory.

A baseline concept, which equips the boattail configuration with a rocket primary cell thruster (presented in Figure 1), has been discussed in the previous sections. For the advanced concept, the aerospike nozzle portion appears as the external nozzle.

Note that in this study, the primary thrust direction is different than that of the installation used in the X-43A [27–28]. The primary thruster in this study is inclined toward the center line of the airframe axis, and therefore, the external nozzle (as an aerospike nozzle) plays a role in the compression wall for the primary engine flow. In contrast to the baseline configuration, a case with an airbreathing primary thruster (i.e., a turbojet and a ramjet) is illustrated in Figure 17(a) as an advanced concept. A fully attached boattail flow, achieved by turning the external flow with a sufficiently long boattail, is presumed for this case.

Figure 17(b) presents another advanced concept that essentially does not possess the boattail portion, only its role in altering the direction of the cell nozzle flow. This configuration possesses the bypass secondary flow. The primary role of the bypass secondary flow is to compensate for the absence of the boattail external flow by creating a virtual boattail flow that flows on the spike surface. The idea behind this configuration that provides the bypass flow is to induce the emergence of an oblique shock wave and form a slip surface between the cell jet primary flow and bypass secondary flow by preventing the cell jet flow from expanding outwardly from the nozzle surface. This concept is different from other thrust augmenting ideas, such as the "base bleed," which is intended to increase the shroud base pressure [19]. The base bleed concept discharged a relatively small amount of low-energy air from the shroud base region, with the goal of destroying the eddy system in the separated base region and substantially increasing the base pressure. Unlike this idea, the concept of providing the bypass flow in this study attempts to utilize the large amount of air flow as the virtual boattail flow for thrust augmentation. For this concept, the airframe-integrated propulsion system is critical in the design of the vehicle because the forebody is part of the inlet, whereas the aft section functions as an exhaust nozzle [23,27]. Additionally, the bypass flow concept combined with an aerospike nozzle is expected to have the advantage of causing less acoustic drag because the plug nozzle system [29] caused less acoustic drag than conventional bell-shaped nozzles.

Because the flight control characteristics, such as the trim of the aircraft, are crucial to the effectiveness of the propulsion system, the flight characteristics with respect to the longitudinal static stability of the aircraft were evaluated by investigating the pitching moment coefficient and the aerodynamic center location. Note that the wing and other attitude control devices were not considered in this analysis. The balance of the pitching moment was evaluated in the steady flight (cruise) condition. This analysis was applied to the configuration presented in Figure 17(a). The x-coordinate of the gravity center ($x_{g.c.}$) was varied from the airframe leading edge to the trailing edge of the spike nozzle in the horizontal direction. The y-coordinate of the gravity center ($y_{g.c.}$) was on the airframe axis, which passes through the airframe leading edge. The pitching moment was calculated around the gravity center by considering the force balance. As a general definition, the positive pitching moment

corresponded to the force acting on the airframe in the nose-up direction, and the negative pitching moment was assumed for the force in the nose-down direction. The force for calculating the positive pitching moment was the pressure force acting on the leading edge surface, and the force for the negative pitching moment was the sum of the pressure force acting on the spike nozzle surface and the cell nozzle thrust. The pressure forces for the forebody were calculated as the average value acting on the entire surface of the leading edge (e.g., $\theta_{LE} = 5^{\circ}$ and $L_{LE} = 10$ m). The distance for calculating the moment was taken as the length between the gravity center and the midpoint of the area where the force was considered. The cruise condition for this analysis was 13 km above sea level, where the atmospheric pressure is approximately 15 kPa. The point where the force balance becomes zero means the force equilibrium, namely the aerodynamic center. For this example case for the baseline configuration (Figure 17(a)) with the aforementioned leading edge configuration, the aerodynamic center was found at 10.9 m from the airframe leading edge, which was located at approximately 9 m upstream of the middle point of the airframe. The locations of the aerodynamic center for other cruise Mach numbers are plotted in Figure 18 along with the calculated case for $\theta_{LE} = 5^{\circ}$ and $L_{LE} = 15$ m. The aerodynamic center was located several meters upstream of the middle point of the airframe for all calculated Mach numbers. Therefore, this proposed configuration is considered to be feasible. A more detailed study on the flight control characteristics will be given in the future work.

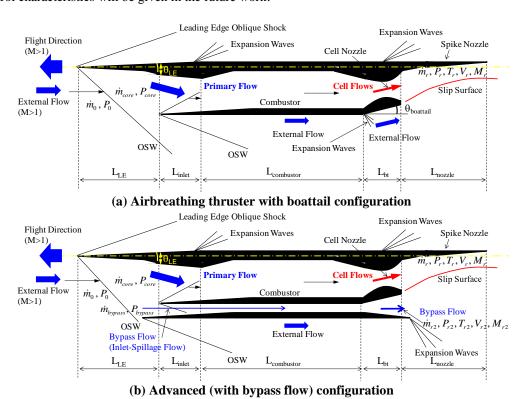


Figure 17 Schematic of airframe/propulsion integrated system for (a) baseline and (b) advanced configurations.

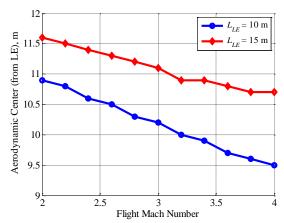


Figure 18 Location of aerodynamic center plotted against cruise Mach number (at altitude of 13 km).

The overall thrust performance prediction model for the configuration presented in Figure 17(a) is given by the summation of the cell nozzle thrust (F_{cell}) and the thrust generated by the spike nozzle ($F_{spike nozzle}$), and expressed by Eqs. (4) to (6). Here, the leading edge drag was omitted to focus on the analysis of the thrust generated by the cell and spike nozzles. Additionally, the term for calculating the base pressures was omitted because the cell and spike nozzle configurations considered here are the 2-D nozzle and full-length spike, respectively.

$$F_{total} = F_{cell} + F_{spikenozzle} \tag{4}$$

$$F_{cell} = \left\{ \dot{m}_r V_r + (P_r - P_0) \cdot A_r \right\} \cdot \cos \theta_{boatsail} - \dot{m}_0 V_0 \tag{5}$$

$$F_{spike nozzle} = \int P_{straight} \cdot \sin \theta_{boattail} \cdot dA_{straight} + \int P_{spike} \cdot \sin(\theta_{boattail} - \theta_{local}) \cdot dA_{spike}$$

$$- \overline{P_{boattail}} \cdot A_{boattail} \cdot \cos \theta_{boattail} - P_b \cdot \int \left\{ \sin \theta_{boattail} \cdot dA_{straight} + \sin(\theta_{boattail} - \theta_{local}) \cdot dA_{spike} + dA_{sb} \right\}$$

$$- \int D_{sf} \cdot \left\{ \sin \theta_{boattail} \cdot dA_{straight} + \sin(\theta_{boattail} - \theta_{local}) \cdot dA_{spike} \right\}$$
(6)

In Eq. (4), the first term in the right-hand side is the thrust by the primary cell thruster and the second term is the thrust generated by the spike nozzle portion, as mentioned above, and each portion is expressed in detail by Eqs. (5) and (6), respectively. In Eq. (5), in the case that a rocket primary thruster is employed, Eq. (5) becomes identical to that of the first term in Eq. (2) by eliminating the ram drag, which appeared as the second term in Eq. (5). When employing the airbreathing primary thruster (i.e., a turbojet and a ramjet), the ram drag term remains. The flow deceleration attributable to the leading edge shock was considered as well.

In the case of the advanced airbreathing concept (shown in Figure 17(b)), a turbojet propulsion is considered up to the flight Mach number range of 3.0; for above this Mach number range, the bypass ramjet propulsion concept is considered. Figure 17 simply illustrates the airframe-integrated propulsion configurations with an emphasis on the difference between the baseline and advanced configurations. Therefore, the mechanical device for switching the propulsion system from turbojet to ramjet, such that shown in a turbine based combined/combination cycle (TBCC) engine, which is usually used to propel a hypersonic airbreathing cruise vehicle [30–32] configuration proposed by Bulman et al. [32], is not considered in detail in this study. As mentioned above, the cell nozzle exit geometry is fixed so that the Mach number at the cell nozzle exit is fixed at 3.5, and the vehicle can be accelerated up to Mach 4.7 for this aerospike nozzle configuration.

The thrust performance prediction model for the cell nozzle and bypass flow for the advanced configuration with a bypass airbreathing primary thruster is given as Eq. (7) by replacing the boattail drag, which appeared as the third term in Eq. (6), with the bypass thrust. In this analysis, the bypass flow is calculated similarly to the general turbofan propulsion. It should be noted that this quasi-one-dimensional analysis simply evaluates the feasibility of the proposed system, which uses the bypass flow to eliminate the boattail drag, because the flowfield is more complex. Therefore, a two- or three-dimensional computational analysis is needed for a more accurate prediction.

$$F_{cell} + F_{bypass} = \dot{m}_r \cdot \left(V_r \cdot \cos \theta_{boattail} - V_{core} + bpr \cdot V_{r2} - bpr \cdot V_{bypass} \right) + \left(P_r - P_0 \right) \cdot A_r \cdot \cos \theta_{boattail} + \left(P_{r2} - P_0 \right) \cdot A_{r2}$$
(7)

Here, the bypass ratio (bpr) is expressed by assuming that the flow rate of fuel is significantly smaller than that of the air flows:

$$bpr = \frac{\dot{m}_{bypass}}{\dot{m}_{core}} = \frac{\dot{m}_{r2}}{\dot{m}_{r}} \tag{8}$$

The thrust equation for the primary cell thruster and bypass flow is rewritten as follows. In the higher supersonic range, a turbojet engine is hypothetically replaced with a ramjet engine because of the drag. The specific heat ratios of the core and bypass flows at the inlet were assumed to be 1.4 because the core and bypass flows are basically ambient air with a value of 1.4. The specific heat ratio for the cell nozzle exhaust gas was presumed to be 1.2. The atmospheric air temperature (T_0) was referred from the U.S. Standard Atmosphere [22] according to the flight altitude. The temperatures for the core flow (T_{core}) and bypass flow (T_{bypass}) at the engine inlet were assumed to be the same as the atmospheric air. The bypass flow Mach number at its exit was assumed to be unity ($M_{r2} = 1.0$); it was assumed that the supersonic external flow incoming to the

bypass flow duct was decelerated, and consequently the flow Mach number at its exit becomes unity. Assuming that $T_{core} = T_{bypass}$ and they can be calculated through the isentropic relation with T_0 and M_{flight} , and that $M_{core} = M_{bypass}$, yields the following relationship:

$$F_{cell+bypass} = \sqrt{\frac{\gamma_r}{R_r T_r}} \cdot P_r \cdot M_r \cdot A_r \cdot \left(M_r \cdot \sqrt{\gamma_r R_r T_r} \cdot \cos \theta_{boattail} + bpr \cdot \sqrt{\frac{2\gamma R T_0}{\gamma + 1}} - (1 + bpr) \cdot M_{flight} \cdot \sqrt{\gamma R T_0} \cdot \left(1 + \frac{\gamma - 1}{2} M_{flight}^2 \right)^{-\frac{1}{2}} \right)$$

$$+ (P_r - P_0) \cdot A_r \cdot \cos \theta_{boattail} + (P_{r2} - P_0) \cdot A_{r2}$$

$$(9)$$

The temperature at the combustor exit was assumed to be 1500 K [24], and that at the cell nozzle exit was calculated by assuming isentropic flow process. The gas constant of the cell nozzle flow was assumed to be 400 J/(kg·K), which is similar to the value for hydrogen combustion with air at 4.0 MPa. The thrust generated by the spike nozzle portion is expressed by Eq. (6), but the boattail drag which is expressed as the third term in Eq. (6) is eliminated by the existence of the bypass flow. The duct areas for the core flow (A_r) and the bypass flow (A_{r2}) were considered to have the same value of 1.0, and the area at the inlet was assumed to be the same as that of the cell nozzle exit. It should be noted that these areas should be optimized for a more accurate performance prediction.

Figure 19 compares the total thrust generated by the primary cell thruster that comprises a rocket or an airbreathing propulsion system (configuration of Figure 17(a)) with the aerospike nozzle portion. The thrust generated by only employing the airbreathing primary thruster is also plotted for comparison. The calculated thrust coefficient is plotted against the flight altitude. The hypothetical flight trajectory was considered as presented in Figure 4. During the cruise phase, the flight Mach number was calculated from the constant dynamic pressure (at 50 kPa) condition. It was assumed that there was no reverse flow due to imbalance between the cell nozzle and bypass flows [23]. The thrust for both configurations generally increases along with the flight altitude, and largely depends on the flight Mach number given for the specific mission phases, as clearly seen as a stepwise change in the lower altitude range. During the cruise phase, which appears for the altitude range over approximately 3.0 km, the thrust gradually decreases as the flight altitude increases. However, the thrust performance in the low altitude range is high. This is mainly caused by the altitude-compensating feature of the aerospike nozzle. At the altitude of approximately 8.5 km, the thrust drops. This is attributable to the fact that the spike nozzle expansion attains its optimal expansion condition at this altitude, and therefore, the spike surface pressure decreased above this altitude, as observed in Figure 8. Then, the thrust keeps decreasing as the altitude increases until approximately 11.1 km, after which the thrust starts to increase for the higher altitude range. This is because of the change in atmospheric temperature; the temperature becomes constant above this altitude as the transition of the atmospheric layer between the troposphere and the stratosphere occurs at

this altitude. Throughout the mission phases from take-off to cruise, the primary thruster configuration of the rocket, which was considered the ideal case, appears to have higher thrust performance than the airbreathing configuration. This is primarily attributed to the ram drag that appeared as the second term in Eq. (5). It is remarkable that the configuration with the aerospike nozzle provides a higher thrust throughout the flight trajectory compared to the case of employing only the airbreathing primary thruster without the aerospike nozzle. Therefore, the proposed configuration, which employs the airbreathing primary thruster with the aerospike nozzle, is beneficial.

Figure 20 presents the total thrust calculated for the baseline and advanced configurations (Figure 17(b)), respectively. The thrust coefficient generally shows a similar trend as observed in Figure 19; the thrust in the cruise phase decreases until reaching the altitude of approximately 11.1 km, then the thrust starts to increase. In the take-off to ascent phases up to the altitude of approximately 6 km, the thrust is higher when utilizing the advanced concept with the accompanying bypass flow than for a pure jet or rocket primary thruster case. However, it is notable that the thrust at the altitude above 11.1 km for the advanced concept with the bypass ratio less than 1.0 shows an increasing trend, while the advanced concepts with the bypass ratio higher than 1.0 show a decreasing trend. With the bypass ratios less than 1.0, the thrust is higher than that of the pure airbreathing case (bpr = 0) up to the altitude of approximately 14 km. In this regard, the advanced configuration with some bypass flow shows a notable increase in the altitude of up to 14 km. With respect to the bypass ratio, a higher thrust performance is expected with the bypass ratio of less than 1.0. This suggests that there would be an optimized point between the bypass ratio and cruise efficiency for a given mission. It should be noted that more realistic configurations for the airbreathing primary thruster cases should be explored by accounting for the inlet design [30,33] simultaneously for optimizing the advanced propulsion system. Additionally, structural installations [1] and weight impacts resulting from flow path modifications should be taken into account.

In summary, by considering the fact that the external flow can aid thrust in off-design flight conditions, as observed in the previous studies [14–15], an advanced propulsion system comprising an airbreathing propulsion system with a bypass flow using a small bypass ratio for the primary cell propulsion system and spike nozzle configuration is beneficial. Additionally, the increase in thrust for the advanced concept can be concluded to add the substantial effect of eliminating the boattail drag. Therefore, the bypass ramjet propulsion system with a linear aerospike nozzle configuration is favorable for the high-speed air transportation system. This conclusion was emphasized by the thrust performance. More specifically, the specific impulse should be considered by accounting for the mission profile and fuel consumption rates. Additionally, airframe structure and weight impacts should also be evaluated.

Certain concerns, such as the required mass flow rate for operating the ramjet propulsion system, inlet design, and drag, should be taken into account; these and other concerns are thoroughly addressed by Bulman et al. [32] for this kind of airframe-integrated propulsion system.

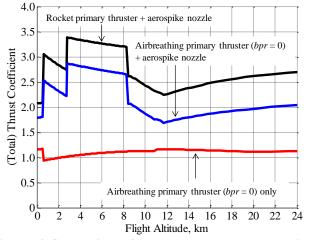


Figure 19 Comparison of total thrust between employing rocket or airbreathing primary thruster along with the aerospike nozzle.

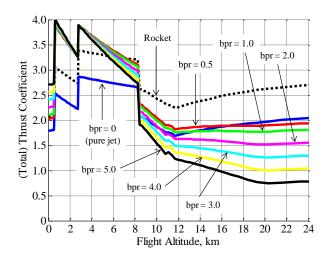


Figure 20 Comparison of total thrust coefficients.

VI. Summary

Clustered linear aerospike nozzles integrated into an airframe and operated under off-design in-flight conditions were evaluated for their thrust performance. The performance evaluation was implemented using the physics-based prediction models established in previous studies. The aerodynamic performance of the nozzle configuration was predicted. Because this study focused on the propulsion system and the airframe portion where the aerospike nozzle was integrated, the aftbody configuration was mainly evaluated in terms of the benefit of the proposed system in increasing thrust performance. The following conclusions were drawn from the performance evaluation of the airframe-integrated linear aerospike nozzle model.

- The boattail-equipped configuration is highly beneficial for increasing thrust because it allows the influence of the external flow to increase the spike surface pressure. Additionally, the boattail angle was found to be the least sensitive to the total thrust.
- 2) The spike length longer than 60% of its full length is beneficial across the complete flight Mach number range considered in this study. Additionally, the spike length was found to be the most important parameter to increase the total thrust.

The following conclusions were drawn from the advanced concept:

The airbreathing propulsion system of the primary thruster followed by the linear aerospike nozzle configuration showed that the employment of the aerospike nozzle portion significantly increases the total thrust over the wide flight range from take-off to supersonic cruise phases. The bypass ramjet propulsion system with bypass ratios less than 1.0 showed benefits for increasing the thrust by compensating for the boattail drag. This notable benefit is clear in low altitude conditions.

In future work, other possible drag components, such as inlet-bypass-bleed drag and inlet-spillage drag, have to be taken into account for more accurate evaluation. Additionally, the reacting exhaust gas should be considered. Further, a more realistic conceptual study, in terms of vehicle operation with trajectory optimization and specific impulse, bypass flow configuration, noise concerns (i.e., jet noise and sonic booms), nozzle acoustics, and weight impacts by altering the flow path to the cruise efficiency has to be conducted.

Acknowledgments

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