# Analysis of swing－by trajectory in three－body problem using B－plane 

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#### Abstract

In general，a swing－by trajectory is designed in a two－body problem in a preliminary design or a case of the trajectory of the spacecraft with high energy．In these cases，the swing－by trajectory is often designed by B－Plane with Patched conics method． However，when the spacecraft has low energy，the approximation of the two－body problem does not hold．Therefore，in this study， the B－Plane used in the two－body problem is extended to the three－body problem and analyze the influence of the secondary body on parameters of the B－Plane of the swing－by trajectories．


## B－plane を用いた三体問題におけるスイングバイ軌道の解析

## 摘要

一般的に，スイングバイ軌道は初期段階での軌道もしくは高エネルギーを持つ宇宙機の軌道では二体問題で軌道設計され る。この場合，スイングバイ軌道は Patched conics 法を用いた B－Planeによって軌道設計が行われることが多い。しかし，低エネルギーを持つ宇宙機の時。二体問題と近似することが出来なくなる。したがって，本研究では，二体問題で用いられ る B－Plane を三体問題に拡張し，スイングバイ軌道のを用いて作られる B－Plane に対していくつかのパラメーターに注目し て 2 天体目の影響を分析する。

## 1．Introduction

Swing－by orbit has been used in many missions and there are many studies on swing－by［1，2］．As examples，Galileo， Voyager1 and Voyager2 utilized the swing－by trajectory．Cur－ rently，DESTINY＋under development by JAXA planes to use a multiple lunar swing－by in order to escape the Earth［3］． Swing－by can change the velocity and the energy of the space－ craft without thrusting．Therefore，it is possible to make the mission term longer and suppress the fuel consumption．

Dynamical models used in trajectory design can be gener－ ally divided into two types，a two－body problem and a multi－ body problem．Although it is desirable to consider the influ－ ence of more celestial bodies，it takes a lot of effort to design the trajectory．For example，the case when the spacecraft has high orbital energy with respect to the Earth is shown in Fig． 1．In this case，since the influence of the gravity other than the Earth is small，the dynamics of the spacecraft can be re－ garded as a the two body problem．The characteristic of de－ signing a trajectory in a two－body problem is that it can be designed analytically．Therefore，the swing－by trajectory is often designed by using B－Plane［4］．B－Plane is a plane per－ pendicular to the speed of the spacecraft toward the swing－by target body，and the parameters on B－Plane has an analytical relationship with the swing－by trajectory．On the other hand， the case when the spacecraft has low orbital energy with re－ spect to the Earth is shown in Fig．2．The approximation of the two－body problem with the Earth does not hold，therefore，the influence of other bodies than be Earth should be considered． In particular，trajectory design using the three－body problem considering the effects of two－celestial bodies is often used． The trajectory in the three－body problem can not be designed analytically．

In the low energy，the analytical solution by B－Plane may
not be valid［5，6］．Therefore，in this study，by designing the B－Plane using the dynamics of the three－body problem，the influence on the parameters constituting the B－Plane is inves－ tigated．Based on the obtained results，we compare the B－ Planes of the two－body problem and the three－body problem， and conduct a physical study on the obtained B－plane．

This paper is organized as follows．Chapter 2 describes how to define B－Plane，and Chapter 3 shows simulation results． Finally，the conclusion is stated．


Fig．1：High Energy Orbit


Fig．2：Low Energy Orbit

## 2. How to define and set B-Plane

## 2. 1 Equation of motion

The mass ration of the primary and secondary bodies are expressed as

$$
\begin{equation*}
\mu_{s e}=\frac{\mu_{e}}{\mu_{e}+\mu_{s}} \tag{1}
\end{equation*}
$$

$\mu_{e}$ and $\mu_{s}$ indicate the gravity constant of the Earth and the Sun, respectively. The equation of motion in two-body problem is expressed as

$$
\begin{align*}
\frac{d^{2} x}{d t^{2}} & =-\frac{\mu_{e}}{R^{3}} x  \tag{2}\\
\frac{d^{2} y}{d t^{2}} & =-\frac{\mu_{e}}{R^{3}} y  \tag{3}\\
\frac{d^{2} z}{d t^{2}} & =-\frac{\mu_{e}}{R^{3}} z \tag{4}
\end{align*}
$$

where

$$
\begin{equation*}
R=\sqrt{x^{2}+y^{2}+z^{2}} \tag{5}
\end{equation*}
$$

The equation of motion in three-body problem is expressed as

$$
\begin{align*}
\frac{d^{2} x}{d t^{2}} & =-\frac{\partial \mathrm{U}}{\partial x}  \tag{6}\\
\frac{d^{2} y}{d t^{2}} & =-\frac{\partial \mathrm{U}}{\partial y}  \tag{7}\\
\frac{d^{2} z}{d t^{2}} & =-\frac{\partial \mathrm{U}}{\partial z} \tag{8}
\end{align*}
$$

where

$$
\begin{align*}
& \overline{\mathrm{U}}=-\frac{1-\mu_{s e}}{r_{1}}-\frac{1-\mu_{s e}}{r_{2}}-\frac{1}{2}\left(1-\mu_{s e}\right) \mu_{s e}  \tag{9}\\
& r_{1}=\sqrt{\left(x+\mu_{s e} \cos t\right)^{2}+\left(y+\mu_{s e} \sin t\right)^{2}+Z^{2}}  \tag{10}\\
& r_{2}=\sqrt{\left(x-\left(1-\mu_{s e}\right) \cos t\right)^{2}+\left(y-\left(1-\mu_{s e}\right) \sin t\right)^{2}+Z^{2}} \tag{11}
\end{align*}
$$

## 2. 2 Definition of each specification

In this study, the Sphere of Influence (SOI) is defined as the distance sufficiently away from the center of the gravity of the swing-by target body, and the distance is calculated by

$$
\begin{equation*}
r_{S O I}=r_{s e}\left(\frac{\mu_{e}}{\mu_{s}}\right)^{\frac{2}{5}} \tag{12}
\end{equation*}
$$

As shown in Fig. 3, the velocity of the spacecraft at $S O I$ assumes the infinite velocity. Here, the infinite velocity is made parallel to the $x$-axis.

## 2. 3 Definition of B-Plane

To generate the B-Plane, first find the position vector $\boldsymbol{B}$ from the center of the Earth when the infinite velocity crosses the B-Plane is calculated. Trajectories are propagated from the perigee to the SOI. At this time, the trajectory is a hyperbolic trajectory obtained analytically. When the asymptotic line of this hyperbolic orbit is parallel to the $x$-axis, which is a straight line connecting the Sun and the Earth, $\boldsymbol{B}$ is a vector whose direction is perpendicular to the asymptote from the center of the Earth. Figure 4 shows this series of flows. Next, when the position vector of the perigee $\boldsymbol{r}_{p}$ is rotated around


Fig. 3: Explanation of values required for swing-by
the x axis, $\boldsymbol{B}$ rotates simultaneously. The surface formed by $\boldsymbol{B}$ is defined as B-Plane. In other words, when the velocity at infinity enters perpendicular to the B-Plane, the trajectory that reaches the perigee of the target is obtained. This is shown in Fig. 5. The angle when the $\boldsymbol{r}_{p}$ is rotated around the x-axis is $\theta_{x}$ and 1 to 360 deg .


Fig. 4: Definition of $\boldsymbol{B}$.


Fig. 5: Definition of B-plane.

Specifically, to find $\boldsymbol{B}$, we make a vector $\boldsymbol{r}_{\text {final }}$ using the viewpoint of $\boldsymbol{v}_{\infty}$ from the center of the Earth. Find the angle $\cos \theta_{s}$ between $\boldsymbol{r}_{\text {final }}$ and $\boldsymbol{v}_{\infty}$ by taking the inner product.

$$
\begin{equation*}
\cos \theta_{s}=\frac{\boldsymbol{v}_{\infty} \cdot \boldsymbol{r}_{\text {final }}}{\left|\boldsymbol{v}_{\infty} \| \boldsymbol{r}_{\text {final }}\right|} \tag{13}
\end{equation*}
$$

Find B.

$$
\begin{equation*}
\boldsymbol{B}=r_{\text {final_point }}-b_{\text {start_point }} \tag{14}
\end{equation*}
$$



Fig. 6: Calculation of B

## 3. Simulation Results

## 3. 1 Simulation for two-body problem

The parameters used in the simulation conditions for the two-body problem are shown in Table 1.

Table 1: Simulation condition values for two-body problems

| Earth's gravity constant $\mu_{e}\left[\mathrm{~m}^{3} / \mathrm{s}^{2}\right]$ | $:$ | $3.986 \times 10^{5}$ |
| :---: | :--- | :---: |
| Sun's gravity constant $\mu_{s}\left[\mathrm{~m}^{3} / \mathrm{s}^{2}\right]$ | $:$ | $1.327 \times 10^{11}$ |
| Earth's revolution radius $r_{s e}[\mathrm{~km}]$ | $:$ | $1.496 \times 10^{8}$ |
| Earth revolution cycle $T_{e}[\mathrm{~s}]$ | $:$ | $3.156 \times 10^{7}$ |
| $r_{S O I}[\mathrm{~km}]$ | $:$ | $9.247 \times 10^{5}$ |
| Peripheral distance $r_{p}[\mathrm{~km}]$ | $:$ | $1.000 \times 10^{4}$ |
| Peripheral speed $v_{p}[\mathrm{~km} / \mathrm{s}]$ | $:$ | 8.929 |
| The angle of the near point vector $\theta_{r_{p}}[\mathrm{deg}]$ | $:$ | $1.192 \times 10^{-3}$ |

The plane spanned by the velocity vector $v_{\infty}$ and $\boldsymbol{B}$ at infinity is defined as the orbital plane. As shown in Fig. 7, let $\theta_{\text {deep }}$ be the angle between $\boldsymbol{B}$ and $\boldsymbol{v}_{\boldsymbol{p}}$ in the orbital plane. Also, as shown in Fig. 8, let $\theta_{\text {cross }}$ be the angle between $\boldsymbol{B}$ and $\boldsymbol{v}_{\boldsymbol{p}}$ around the x axis. The simulation results using $\boldsymbol{B}$ and $\boldsymbol{r}_{p}$ are


Fig. 7: Definition of find $\theta_{\text {deep }}$
shown in Fig. 11 and Fig. 12. $\theta_{\text {deep }}$ was constant at -5.969 deg and $\theta_{\text {cross }}$ was almost constant at 0.00 deg , where, the radius of the circle indicates the size of $\boldsymbol{B}$, and the color bar indicates


Fig. 8: Definition of find $\theta_{\text {cross }}$
the size of each angle.


Fig. 9: Angle $\theta_{\text {deep }}$ between $\boldsymbol{B}$ and $\boldsymbol{v}_{\boldsymbol{p}}$ in the orbital plane


Fig. 10: The angle $\theta_{\text {cross }}$ between $\boldsymbol{B}$ and $\boldsymbol{v}_{\boldsymbol{p}}$ around the x axis

## 3. 2 Simulation for three-body problem

Table 2 shows the parameters used in the simulation conditions for the three-body problem. This time, in the three-body problem of the sun-earth system, the Earth's revolution cycle is $2 \pi$ and the earth's revolution radius is 1 .

Table 2: Simulation condition values for three-body problems

| Non-dimensional gravity constant $\mu_{s e}[-]$ | $:$ | $3.003 \times 10^{-6}$ |
| :---: | :--- | :---: |
| Earth's revolution radius $r_{s e}[-]$ | $:$ | 1.000 |
| Earth revolution cycle $T_{e}[-]$ | $:$ | $2 \pi$ |
| $r_{S O I}[-]$ | $:$ | $6.181 \times 10^{-3}$ |
| Peripheral distance $r_{p}[-]$ | $:$ | $6.6845 \times 10^{-5}$ |
| Peripheral speed $v_{p}[-]$ | $:$ | 0.2998 |
| The angle of the near point vector $\theta_{r_{p}}[\mathrm{deg}]$ | $:$ | $1.192 \times 10^{-3}$ |

The vector $\boldsymbol{B}$ and B-Plane in the same way as in the two-
body simulation. Based on this $\boldsymbol{B}$, the values of $\theta_{\text {deep }}$ and $\theta_{\text {cross }}$ are found. However, the radius of the circle indicates


Fig. 11: Angle $\theta_{\text {deep }}$ between $\boldsymbol{B}$ and $\boldsymbol{v}_{\boldsymbol{p}}$ in the orbital plane


Fig. 12: The angle $\theta_{\text {cross }}$ between $\boldsymbol{B}$ and $\boldsymbol{v}_{\boldsymbol{p}}$ around the x axis
the size of $\boldsymbol{B}$, and the color bar indicates the size of each angle.

It was found that $\theta_{\text {deep }}$ takes the maximum value when $\theta_{x}=$ 180 deg , and takes the minimum value when $\theta_{x}=360 \mathrm{deg}$. It was also found that $\theta_{\text {cross }}$ takes the maximum value when $\theta_{x}=91 \mathrm{deg}$, and takes the minimum value when $\theta_{x}=269 \mathrm{deg}$.

## 4. Conclusion

In this study, we investigated the effect of the B-Plane used in the trajectory design for the two-body problem on the parameters constituting the B-Plane when the three-body problem was designed with the dynamics. It was confirmed that the B-Plane changed under the influence of the Sun. Future work includes the analysis of swing-by trajectories with different energies to give more physical insights into trajectory design.

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